

# Influence of gate capacitance on CNTFET performance using Monte Carlo simulation

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# Plan

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- Carbon nanotubes: high potential material
- Carbon nanotubes: Model for calculations (Monte Carlo)
- Evaluation of Mean Free Path
- Carbon nanotube field effect transistor simulation

# Potentialities of Carbon nanotubes

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## Charge transport and confinement

- ✓ Charge transport is quasi-ballistic: few interactions
- ✓ Charge confinement (1D): good control of electrostatics

## No dangling bonds

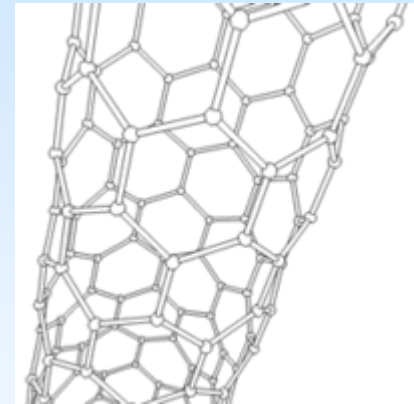
- ✓ Possibility to use High K material for oxide

## Band structure « quasi » symmetric

- ✓ Transport properties identical for electron and hole (same  $m^*$ )

## Carbon nanotubes can be either metallic as well as semi-conducting

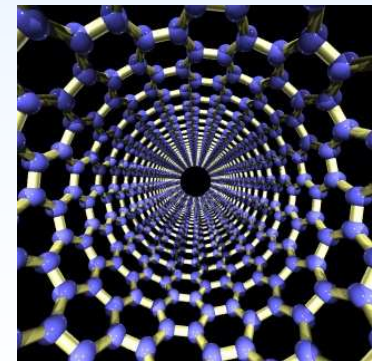
- ✓ All-nanotube-based electronic is possible



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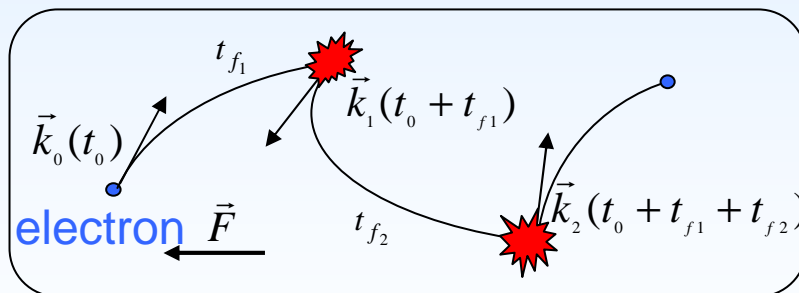
# Particle Monte-Carlo method

Study of carrier transport in this work: solving the *Boltzmann* equation by the Monte Carlo method

$$\frac{\partial f}{\partial t} = -\vec{v} \cdot \vec{\nabla}_r f - \frac{\vec{F}}{\hbar} \cdot \vec{\nabla}_k f + \left. \frac{\partial f}{\partial t} \right|_{\text{collisions}}$$

Statistical solution :by **Monte Carlo method**

- $f(\vec{r}, \vec{k}, t)$  → assembly of individual particles
- 1 particle  $\Leftrightarrow \vec{r}(t), \vec{k}(t)$
- $N$  particles allow us to reconstruct  $f(\vec{r}, \vec{k}, t)$



Carrier trajectories:

Succession of free flight and scattering events

scattering rates  $\lambda_i(\vec{k})$



random selection of

(Monte Carlo algorithm)

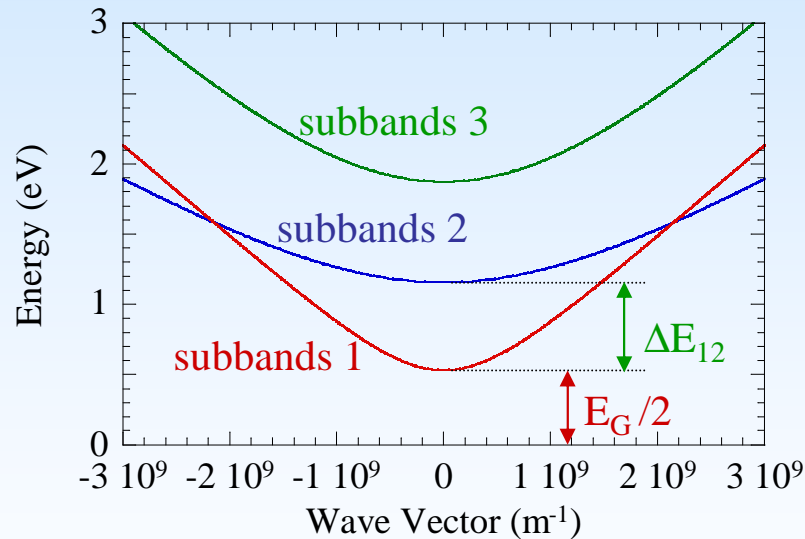


- time of free flights  $t_f$
- type of scattering  $i$
- effect of scattering  $(\Delta E, \theta)$

# Band structure and phonons

Energy dispersion calculated by:

- Tight-binding
- Zone-folding



Analytical approximation of the three first subbands

$$\frac{\hbar^2 k_z^2}{2m_p} = [E_p - E_p^m] \left[ 1 + \alpha_p (E_p - E_p^m) \right]$$

Phonon dispersion curves calculated by zone folding method

## Approximation

➤ Longitudinal acoustic and optical modes

are considered as **dominant** like in graphene

Analytical approximation of dispersion curves  
[Pennington, *Phys. Rev. B* 68, 045426 (2003)]

+ RBM phonon ( $E_{\text{RBM}} \approx \text{Cst.}/d_t$ )

[M. Machon et al., *Phys. Rev. B* 71, 035416, (2005)]

Ex: Tube index  $n=19$ ,  $E_{\text{RBM}}=19 \text{ meV}$

# Scattering Events Calculation

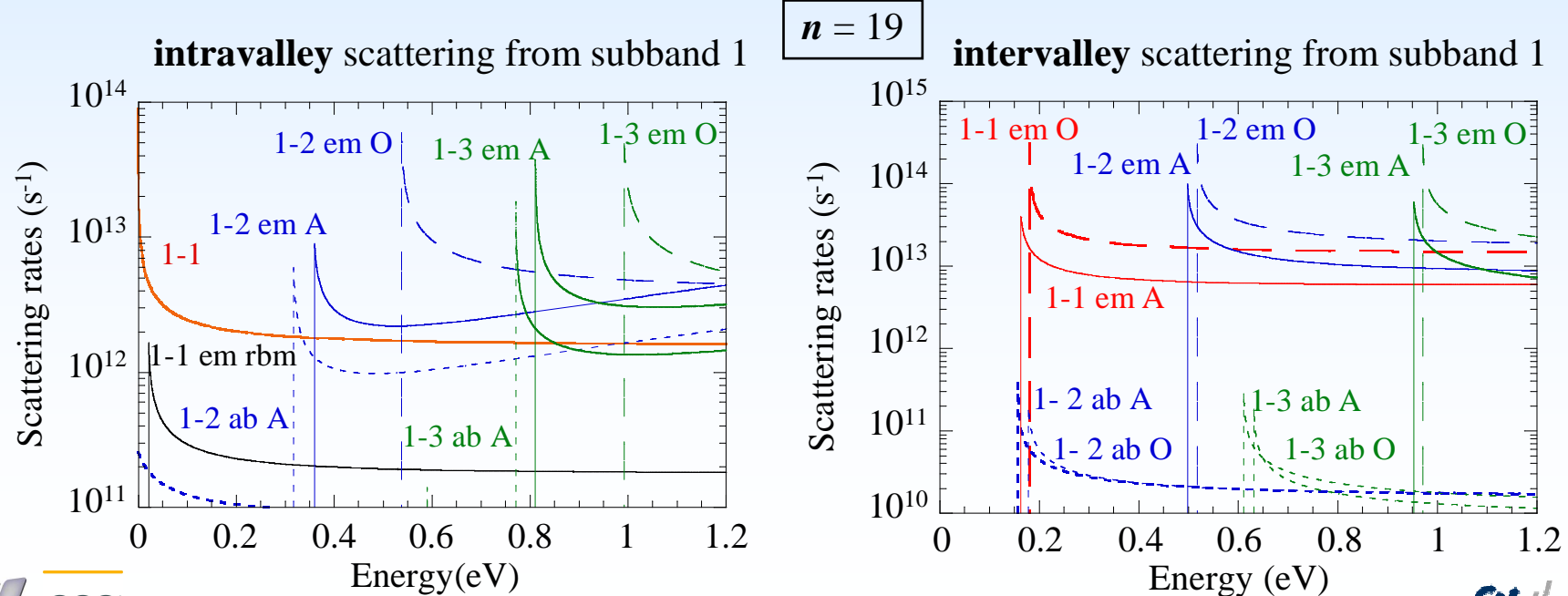
- **Scattering rates** : 1st order Perturbation Theory by deformation potential model

with:  $D_{aco} = 9 \text{ eV}$  [L. Yang, M.P. Anantram et al. Phys. Rev. B **60**, 13874 (1999)]

$D_{RBM} = 0.65 \text{ eV/cm}$  (for  $n = 19$ ) [M. Machon et al., PRB 71, 035416 (2005)]

☞ **Intrasubband acoustic scattering**  $\longrightarrow$  *Elastic* process

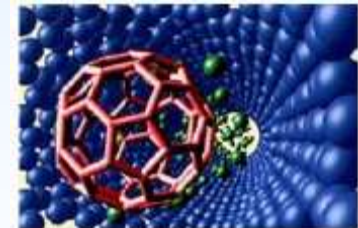
☞ **Intersubband transition**  $\longrightarrow$  *Inelastic* process ( $E = \hbar\omega$ )



# Plan

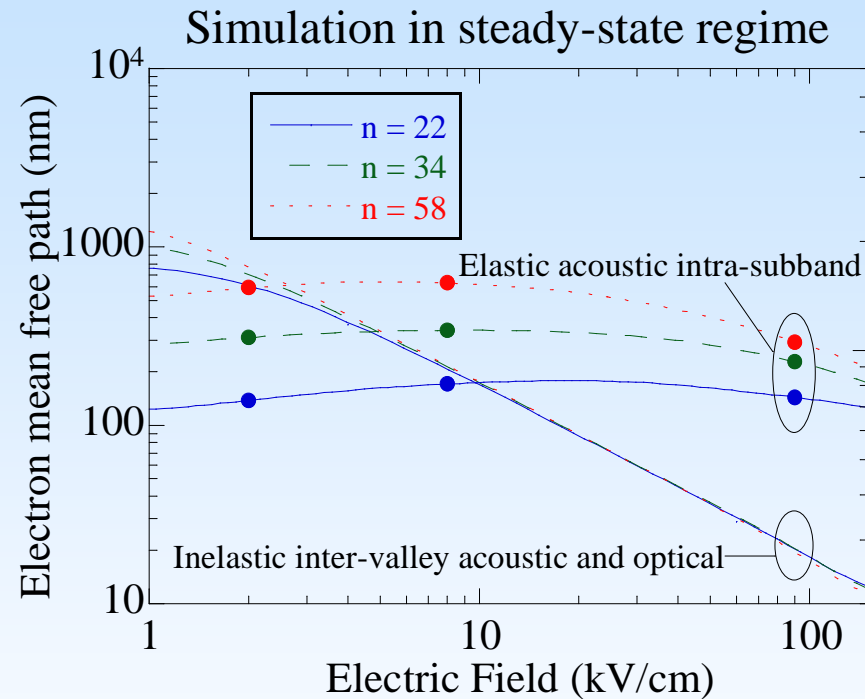
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# Mean Free Path



The mean free path depends on diameter and electric field :

☞  $l_{\text{MFP}} = 100\text{-}600$  nm for elastic scattering

☞  $l_{\text{MFP}} = 1000\text{-}20$  nm for inelastic scattering

Agreement with experimental works (L. field: 300-500 nm H. Field: 10-100nm)

➔ Potential of quasi-ballistic transport in the low-field regime

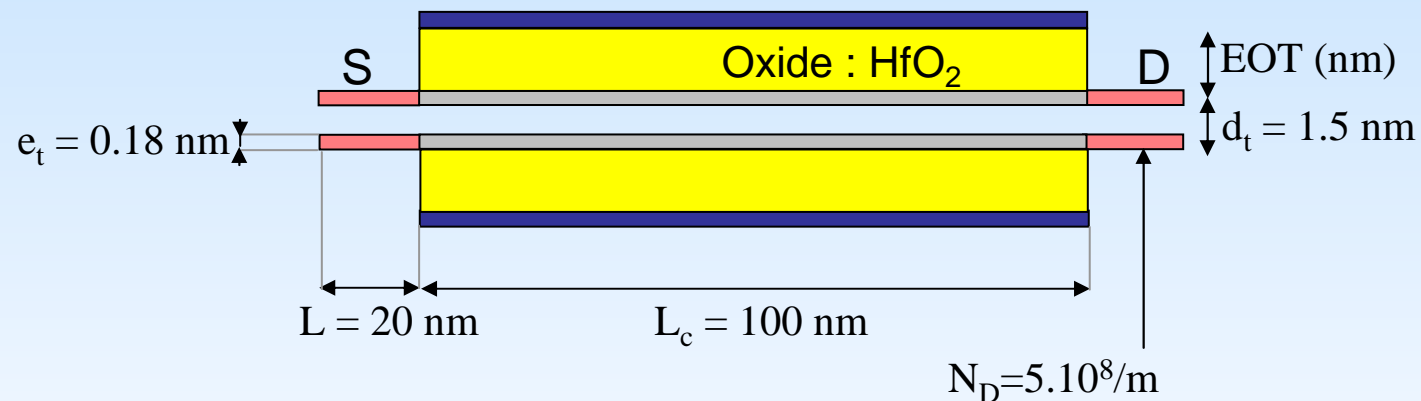
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# The modelled transistor

Coaxially gated carbon nanotube transistor with heavily-doped source drain extensions CNTFET

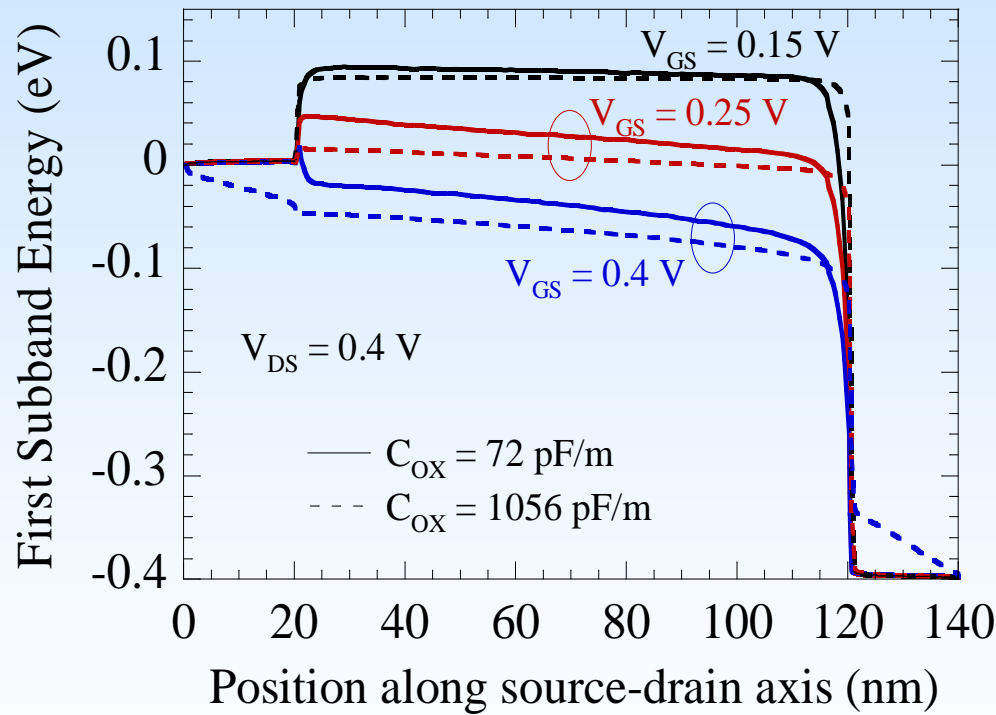


## Characteristics:

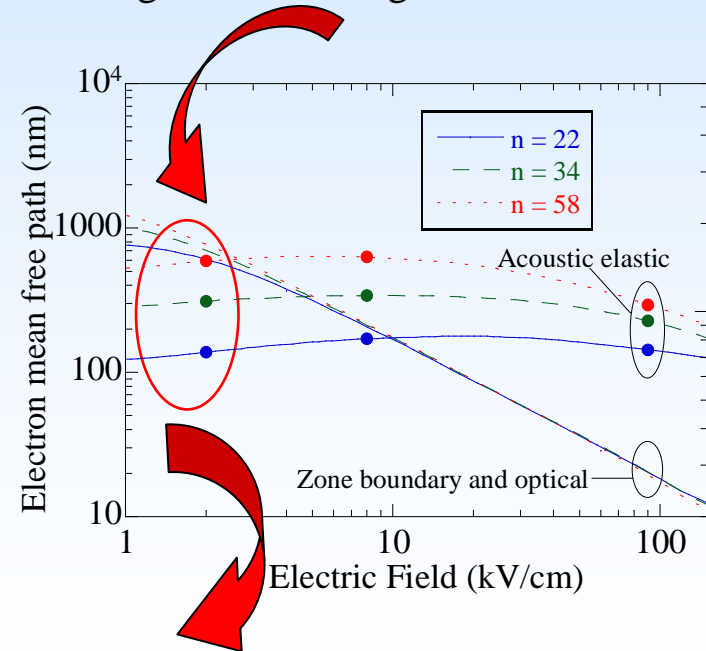
- Zigzag tube (19,0) -  $d_t = 1.5$  nm -  $E_G = 0.55$  eV
- Cylindrically gated
- Ohmic Contacts (heavily n-doped)
- Intrinsic channel
- High K gate insulator HfO<sub>2</sub> ( $EOT = 16, 1.4, 0.4, 0.1$  nm -  $C_{OX} = 72, 210, 560, 1060$  pF/m)

# Conduction band

First subband energy profile



Weak field in the channel due to strong electrostatic gate control



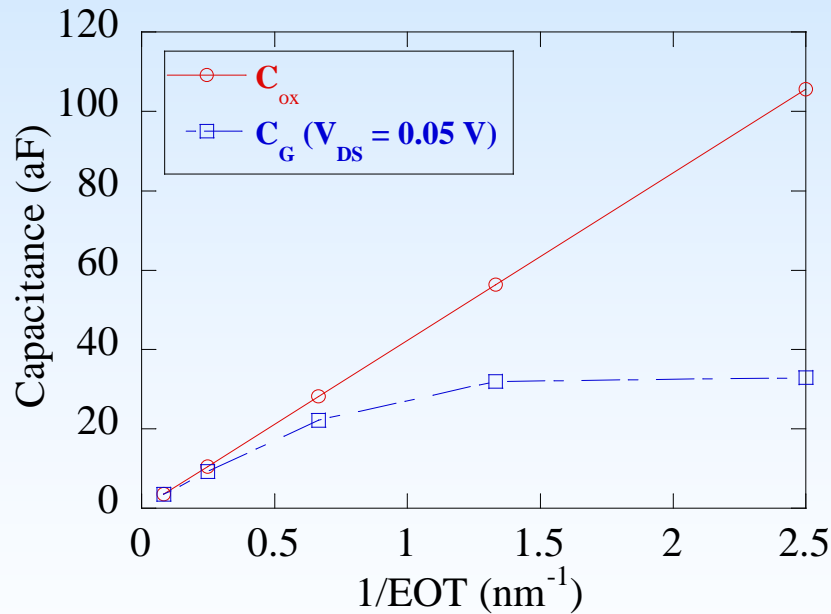
Weak field in the channel  $\rightarrow$  quasi ballistic transport

# Capacitive effects

Weak Field  $\leftrightarrow$  1D low density of states  $\rightarrow$  Channel potential is fixed

Determination of  $C_G$  for low  $V_{DS}$ :  $C_G = \frac{dQ}{dV_{GS}}$       $C_{ox} = \frac{2\pi\epsilon_0\epsilon_r}{\ln\left(\frac{e_{ox} + r_{dt}}{r_{dt}}\right)}$

Evolution of  $C_{ox}$  and  $C_G$  as a function of  $1/EOT$



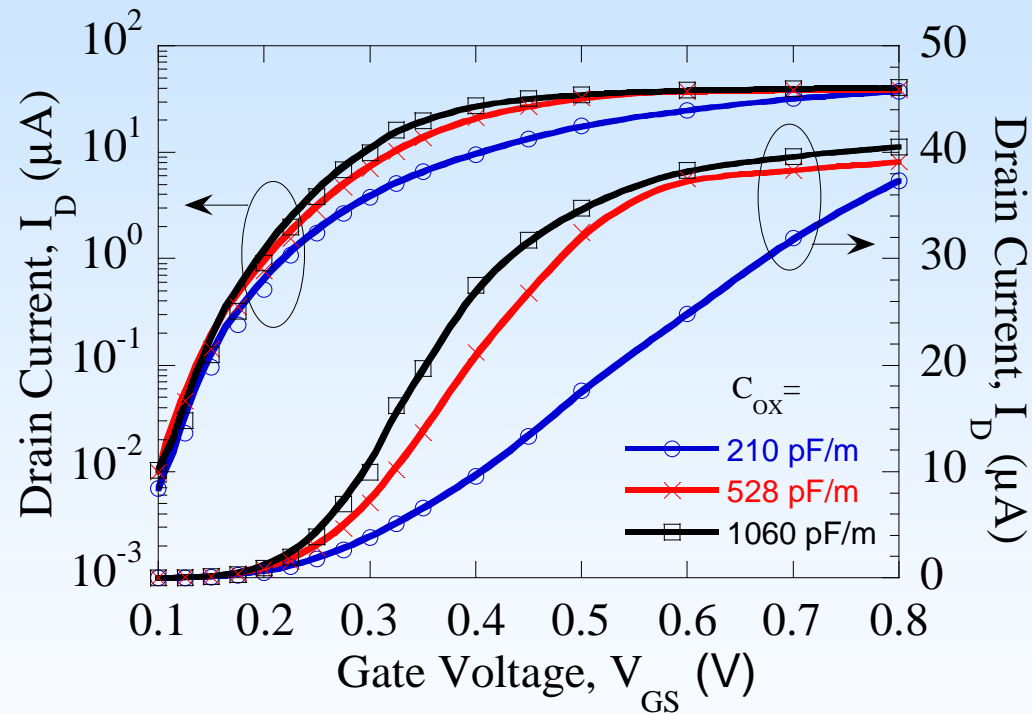
Reminder :  $\frac{1}{C_G} = \frac{1}{C_{ox}} + \frac{1}{C_Q}$

$C_G$  leads towards  $C_Q$  for  $C_{ox} \gg C_Q$

$\approx C_G$  tends towards 4pF/cm

For  $C_{ox} \gg C_Q \rightarrow$  quantum capacitance limit

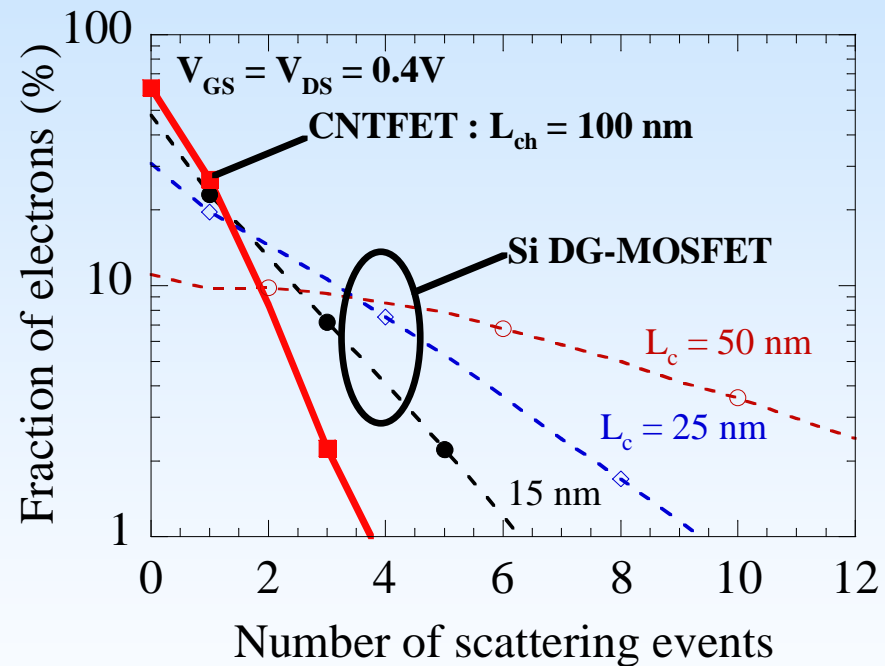
# $I_D$ - $V_{GS}$ Characteristics



Expected result :  $I_D$  increases with oxide capacitance  $C_{ox}$   
but limited enhancement for high  $C_{ox}$  values

# Evaluation of Ballistic Transport

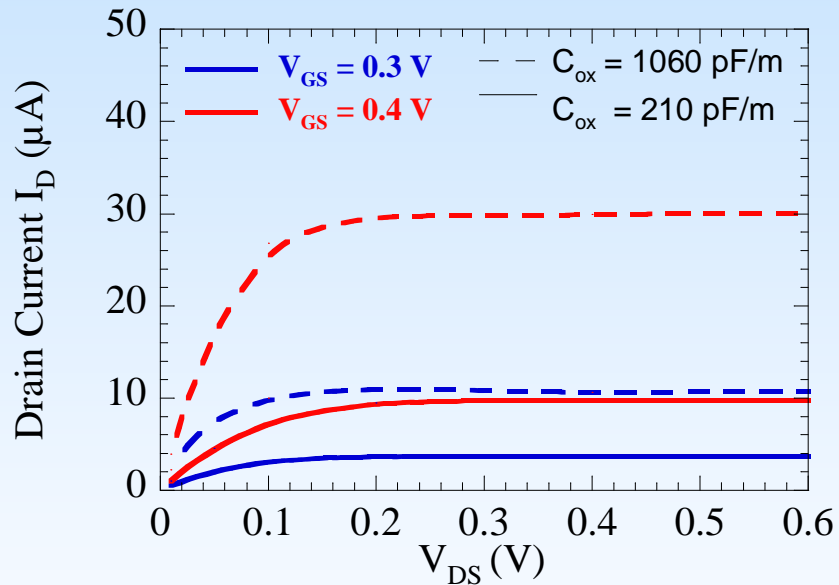
Unstrained Si. DG-MOSFET vs c-CNTFET



**CNTFET : 70% of ballistic electrons at the drain end for  $L_{ch} = 100\text{nm}$**

**Unstrained Si DG : 50% of ballistic electrons pour  $L_{ch} = 15 \text{ nm}$**

# $I_D$ - $V_{DS}$ characteristics



Device parameters extracted from simulation ( $I_{off}=0.1nA$ )

$C_{ox}$ (pF/m) /EOT(nm)	210/ 1.4	560/ 0.4	1060/ 0.1	Expe.* 1.5
$S$ (mV/dec)	70	65	60	70
$I_{ON}$ ( $\mu A$ )	24	40	44	5
$g_m$ ( $\mu S$ )	32	62	100	20

$V_{DD} = 0.4V$

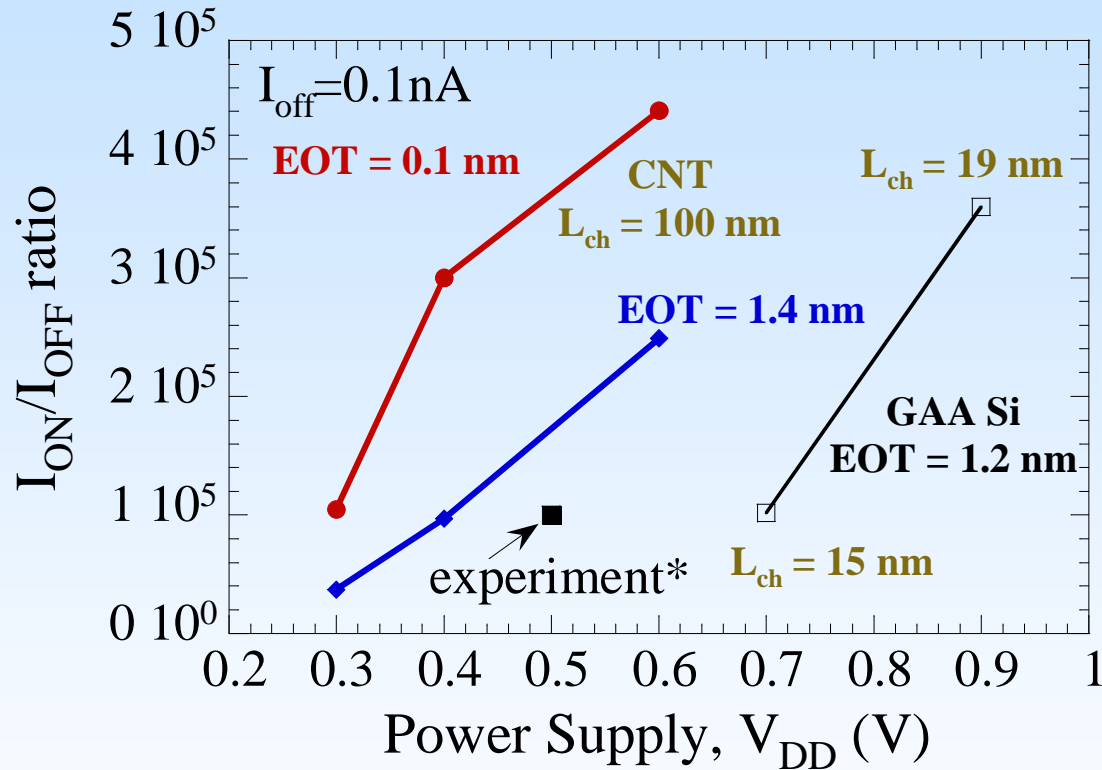
$I_{ON} \approx 10 - 30 \mu A$ .  $I_{ON}/I_{OFF} \approx 510^4 - 510^5$

High values of  $S$  and  $g_m$  are obtained  
Acceptable agreement with experiments

\* *H. Dai, Nano Letter, vol.5, 345, (2005) :  $L=80nm$ ,  $EOT=1.5nm$*



# $I_{ON}/I_{OFF}$ ratio ( $V_{DD}$ )



Performances CNT for  $L_C=100\text{nm}$  and  $V_{DD}<0.4\text{V}$



Performances Si for  $L_C=15\text{nm}$  and  $V_{DD}>0.7\text{V}$

\* H. Dai, Nano Letter, vol.5, 345, (2005) :  $L=80\text{nm}$ ,  $EOT=1.5\text{nm}$

# Conclusion

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**This Monte Carlo simulation of a CNTFET shows excellent performances:**

- Excellent electrostatic gate control (quantum capacitance regime)
- High fraction of ballistic electrons in the transistor channel
- High performances at relatively low power supply

